

What is Q?

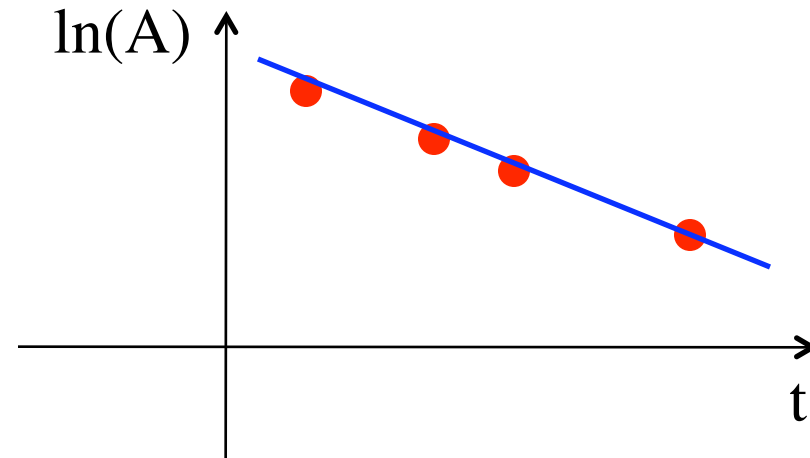
Interpretation 1: Suppose A_0 represents wave *amplitudes*, then

$$A = A_0 e^{-bt} = A_0 e^{-\omega_0 t / (2Q)}$$

$$\ln(A) = \ln(A_0) - \left[\frac{\omega_0}{2Q} \right] t$$

intercept

slope



Interpretation 2: Suppose u represents *displacement*, then

$$u(t) = A_0 e^{i(a+ib)t} = A_0 e^{i\omega_0 t} e^{-\omega_0 t / (2Q)}$$

$$a = \omega_0 \sqrt{1 - 1/4Q^2} \quad (\text{real}) \longrightarrow \omega = \text{“modified” or “instantaneous freq”}$$

$$b = \frac{\omega_0}{2Q} \quad (\text{imaginary})$$

Suppose: small attenuation, the $\omega \approx \omega_0$ $b \approx \frac{\omega}{2Q}$

We can define $b = \omega^*$, where $\omega^* \rightarrow 0$ as Q increases (imaginary freq due to attenuation),

$$\longrightarrow \omega^* = \frac{\omega}{2Q} \quad \longrightarrow Q^{-1} = 2\omega^* / \omega$$

Relation with velocity:

$$c + ic^* = \frac{\omega}{k} + i \frac{\omega^*}{k} = \frac{\omega}{k} + i \frac{Q^{-1}}{2k}$$

\longrightarrow Imaginary velocity due to attenuation

$$c^* = \frac{\omega}{2k} Q^{-1} \Rightarrow Q^{-1} = 2 \frac{c^*}{c}$$

So Q is a quantity that defines the relationship between real and imaginary frequency (or velocity) under the influence of attenuation.

Interpretation 3: Q is the number of cycles the oscillations take to decay to a certain amplitude level. $n = t/T = t(\omega/2\pi)$

$$\text{if } Q \rightarrow \infty \quad \text{then } n \approx t(\omega_0/2\pi) \rightarrow t_n = \frac{n \cdot 2\pi}{\omega_0}$$

So amplitude at time t_n (after n cycles)

$$A(t_n) \approx A_0 e^{-\omega_0 t_n / 2Q} = A_0 e^{-\omega_0 n 2\pi / (\omega_0 2Q)} = A_0 e^{-n\pi / Q}$$

$$\text{if } n = Q, \text{ then } A = A_0 e^{-\pi} \approx 0.04 A_0$$

Attenuation and Physical Dispersion (continued...)

Different interpretations of Q (quality factor):

- (1) As a damping term $Q = \omega_0/\gamma$
- (2) As a fraction between imaginary and real frequencies (or imaginary velocity to real velocity)

$$Q^{-1} = \frac{2\omega^*}{\omega} \quad \text{or} \quad Q^{-1} = \frac{2c^*}{c}$$

- (3) As the number of cycles for a wave to decay to a certain amplitude. If $n = Q$, then

$$A(t_n) = A_0 e^{-\omega_0 t / 2Q} = A_0 e^{-(\omega_0 / 2Q)(n * 2\pi / \omega_0)} \approx 0.04 A_0$$

- (4) Connection with t^* (for body wave).

$$t^* = \int_{path} \frac{dt}{Q(\mathbf{r})} \approx \sum_{i=1}^N \frac{\Delta t_i}{Q_i} \quad (\text{N layers, } \mathbf{r} = \text{location})$$

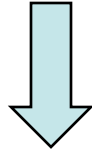
1 number that describes Q structure of several layers

- (5) Energy formula $\frac{1}{Q(\omega)} = -\frac{\Delta E}{2\pi E}$ ($-\Delta E$ = energy loss per cycle)

Effects of Q (assume the SAME Q value)

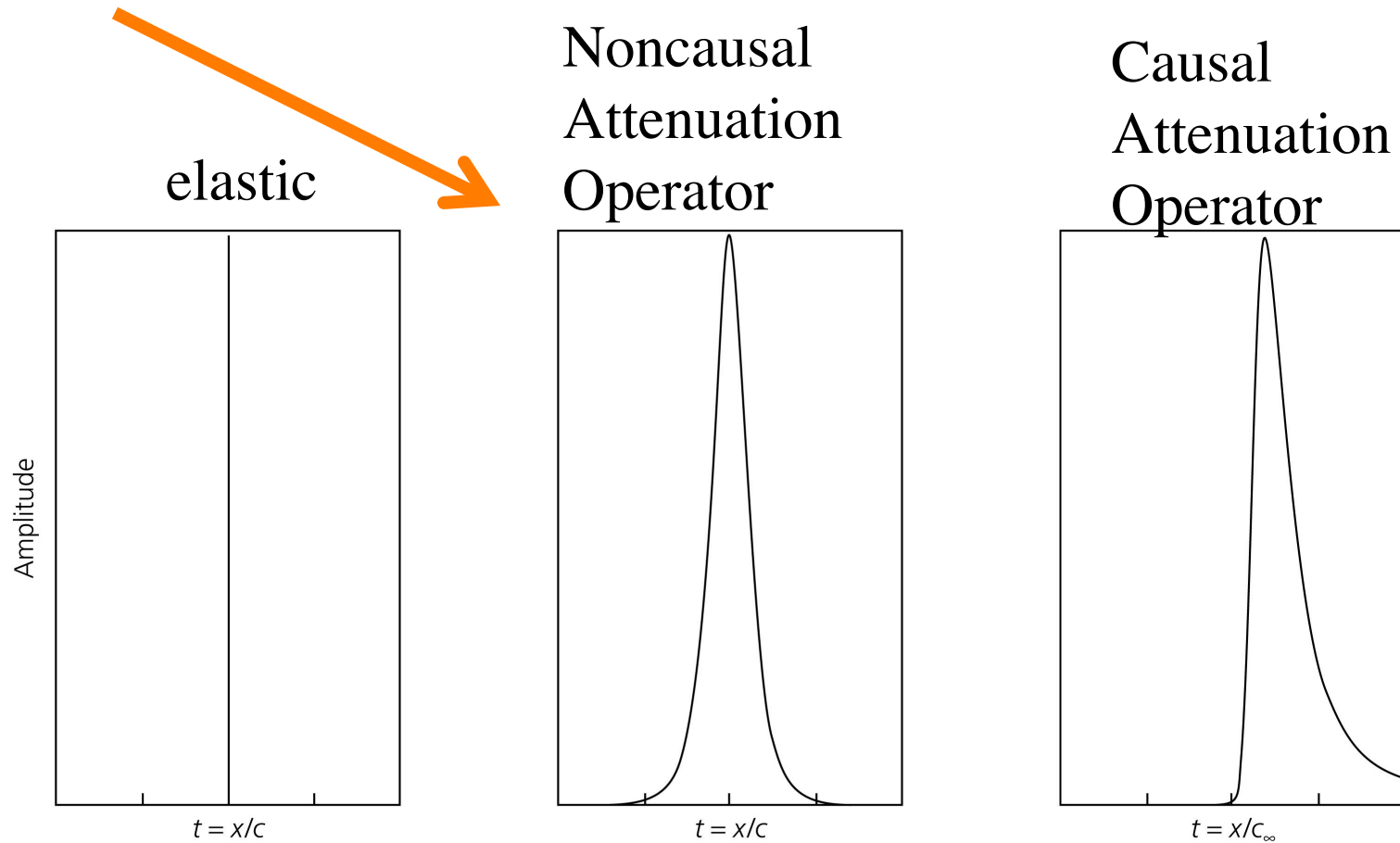
$$(1) \quad A(\omega) = A_0(\omega)e^{-\omega_0 t / 2Q} \approx A_0(\omega)e^{-\omega x / (2Qv)}$$

dependencies



- { large distance x ----- > more amplitude decay
- { large velocity v ----- > the less amplitude decay
- { large frequency ω ----- > more amplitude decay

Physical Dispersion



Observation: pulse is “spread out” which means dispersive (different frequencies arrive at different times!)

Problem: envelope of the function is nonzero before $t=x/c$ (it is like receiving earthquake energy before the rupture, not physical!) 5

How to make this process causal?

*One of the often-cited solutions: **Azimi's Attenuation Law:***

$$c(\omega) = c_0 \left[1 + \frac{1}{\pi Q} \ln \left(\frac{\omega}{\omega_0} \right) \right] \quad c_0 = \text{reference speed for frequency } \omega_0$$

if $Q = \infty$, then $c = c_0$ (no dispersion)

if $\omega = \omega_0$, then $c = c_0$ (no dispersion)

if $\omega = \infty$, then $c \approx c_0$ (means high freq arrives first)

Where does it come from?

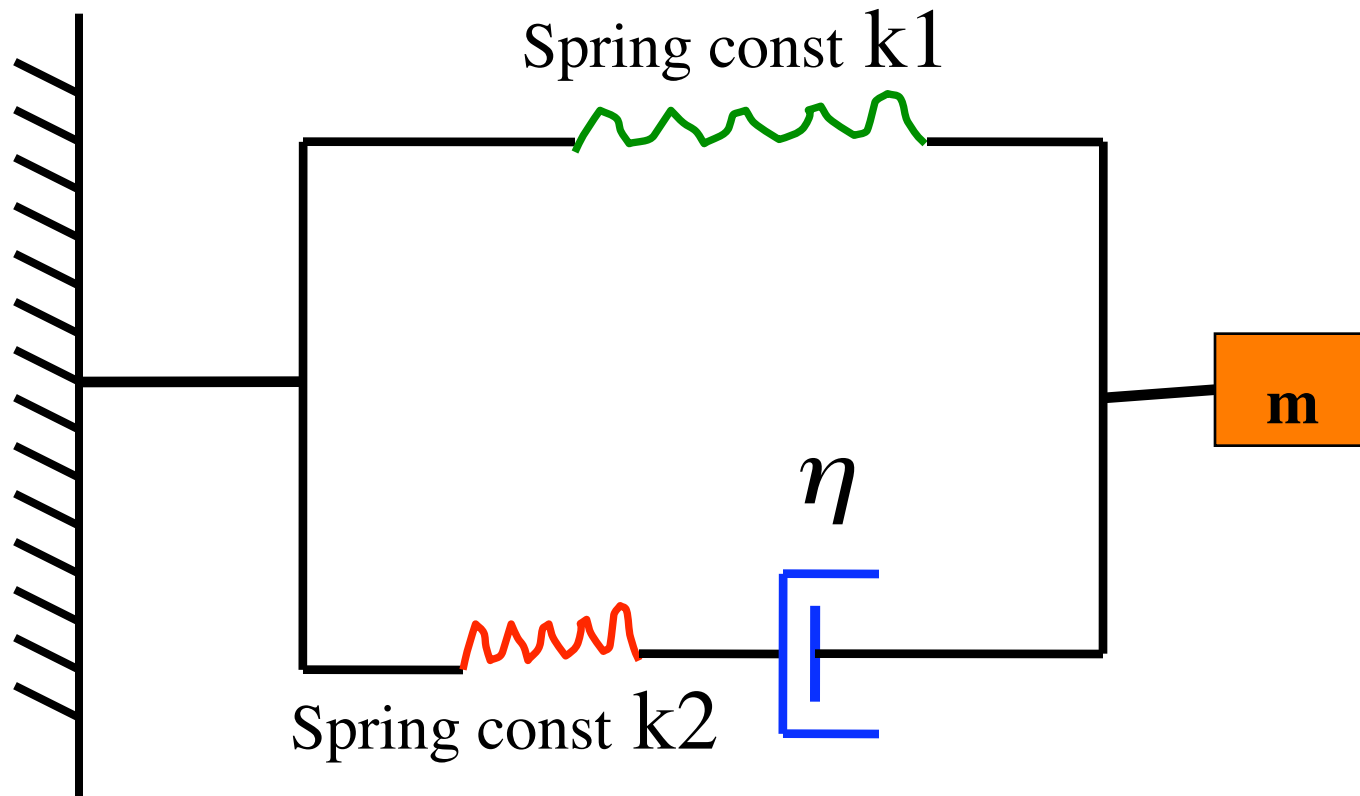
Answer: Derived under the following **causality condition:**

$u(x, t) = 0$ for all $t < x/c(\infty)$, where $c(\infty)$ is the highest (infinite) frequency that arrives first.

Physical Models of Anelasticity

In a nutshell, Earth is composed of lots of **Viscoelastic (or Standard Linear) Solids**

Standard Linear Solid (or Viscoelastic Solid)



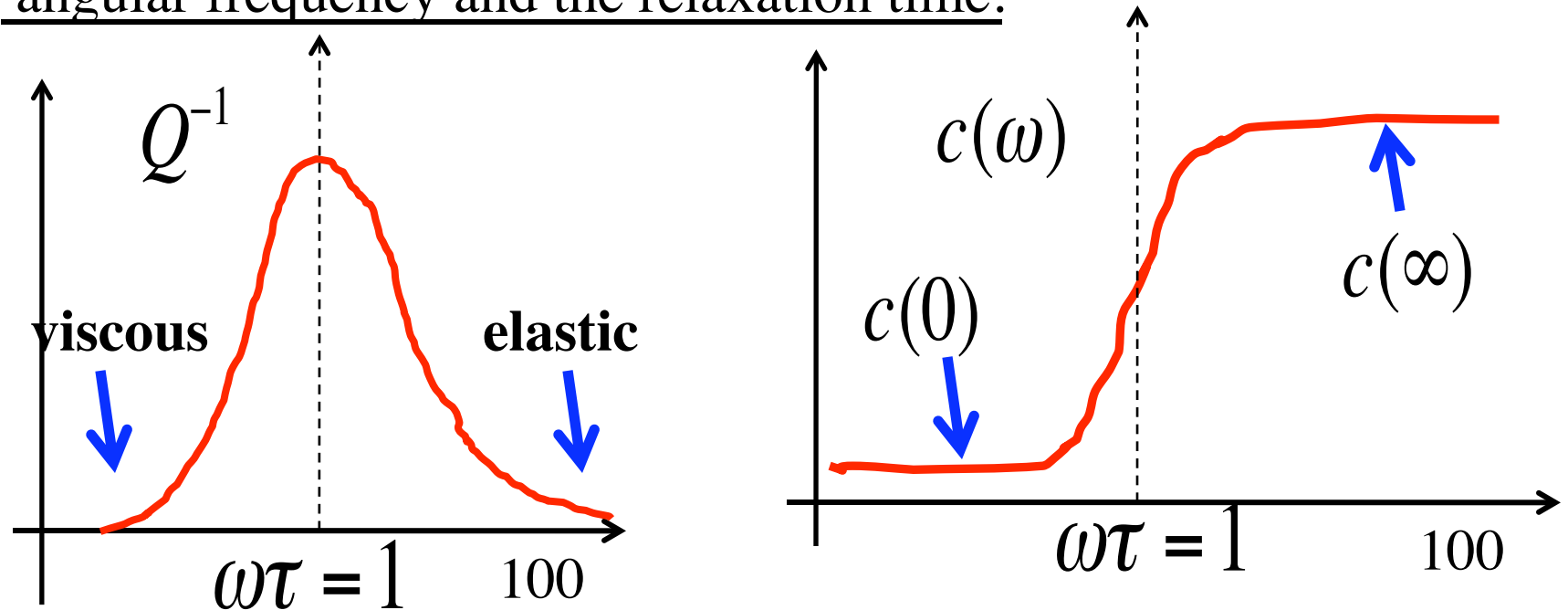
Specification: (1) consist of two springs and a dashpot
(2) η = viscosity of fluid inside dashpot

Governing Equation (stress): $\sigma(t) = k_1 H(t) + k_2 e^{-t/\tau}$

where $H(t)$ = step (or Heavyside function)

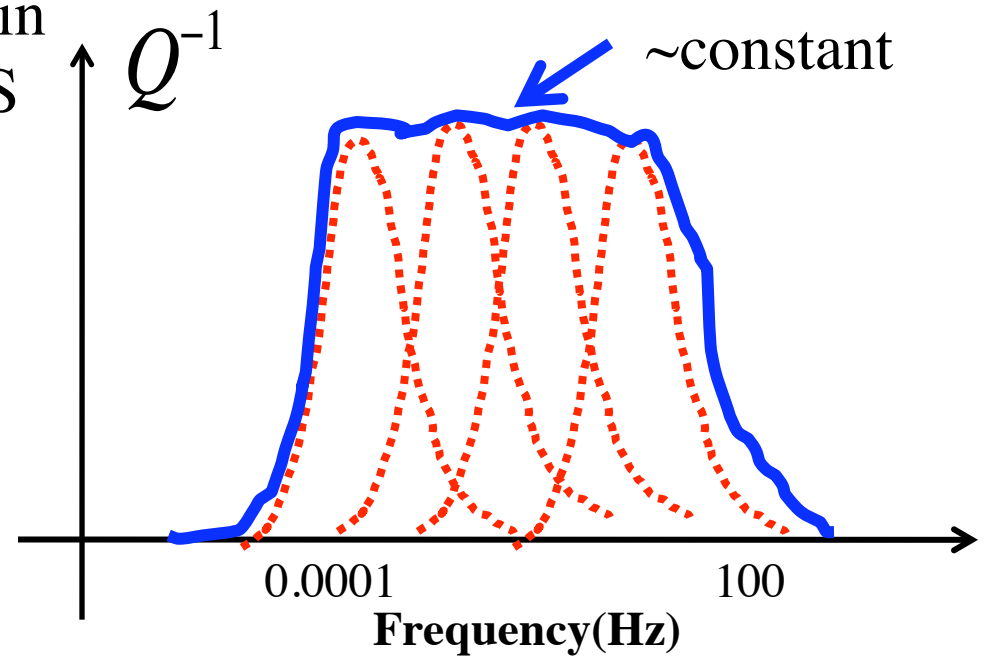
and $\tau = \eta/k_2$ (relaxation time)

The response to a harmonic wave depends on the product of the angular frequency and the relaxation time.



The left-hand figure is the absorption function. The absorption is small at both very small and very large frequencies. It is the max at $\omega\tau = 1$!

A given polycrystalline material in the Earth is formed of many SLS superimposed together. So the final Frequency dependent Q is constant for many seismic frequencies ----→

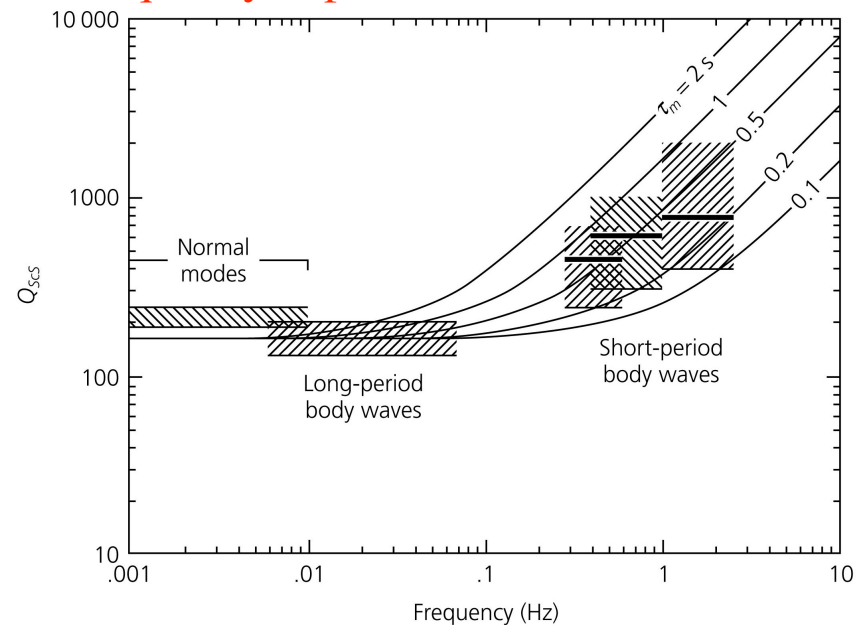


Question: Is the fact that high frequencies are attenuated more a contradiction to this flat Q observation??

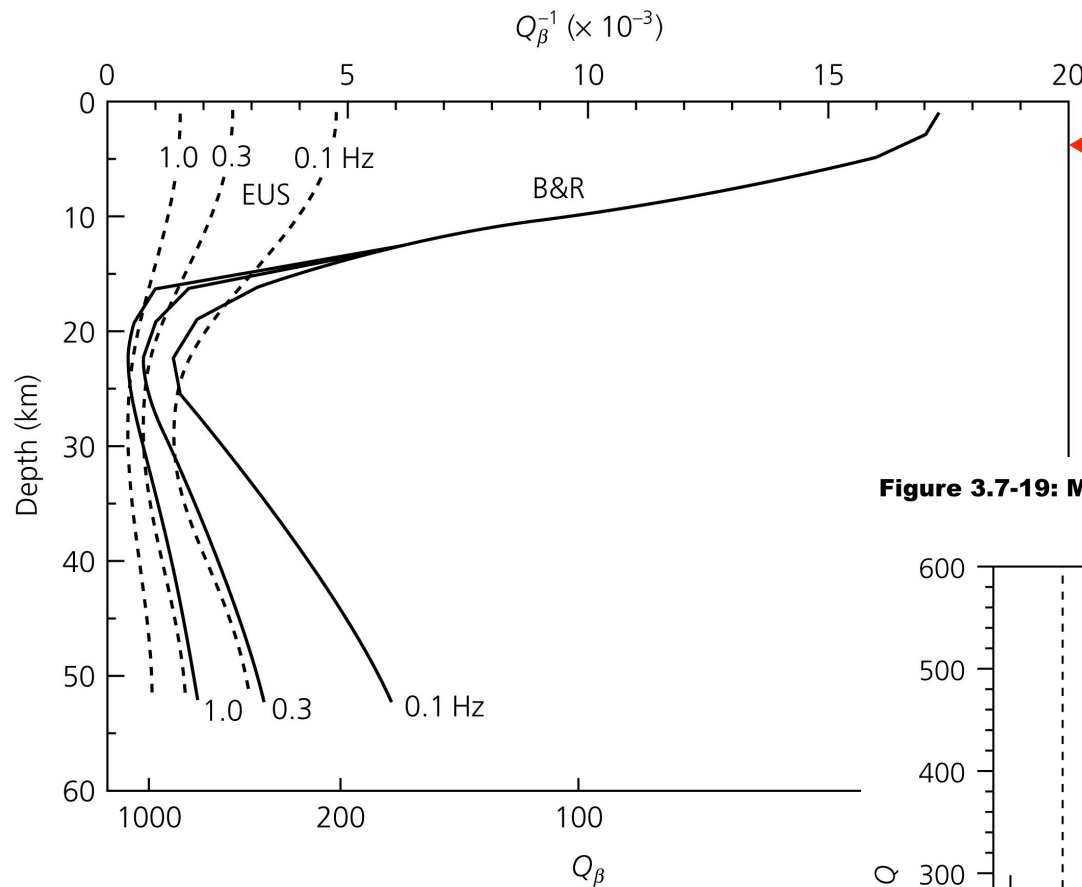
Answer: No. High frequencies are attenuated more due to the following equation that works with the same Q . So it is really “frequency dependent amplitude”, NOT “frequency dependent Q ”.

$$A(\omega) \approx A_0(\omega)e^{-\omega x/(2Qv)}$$

Frequency dependence of mantle attenuation

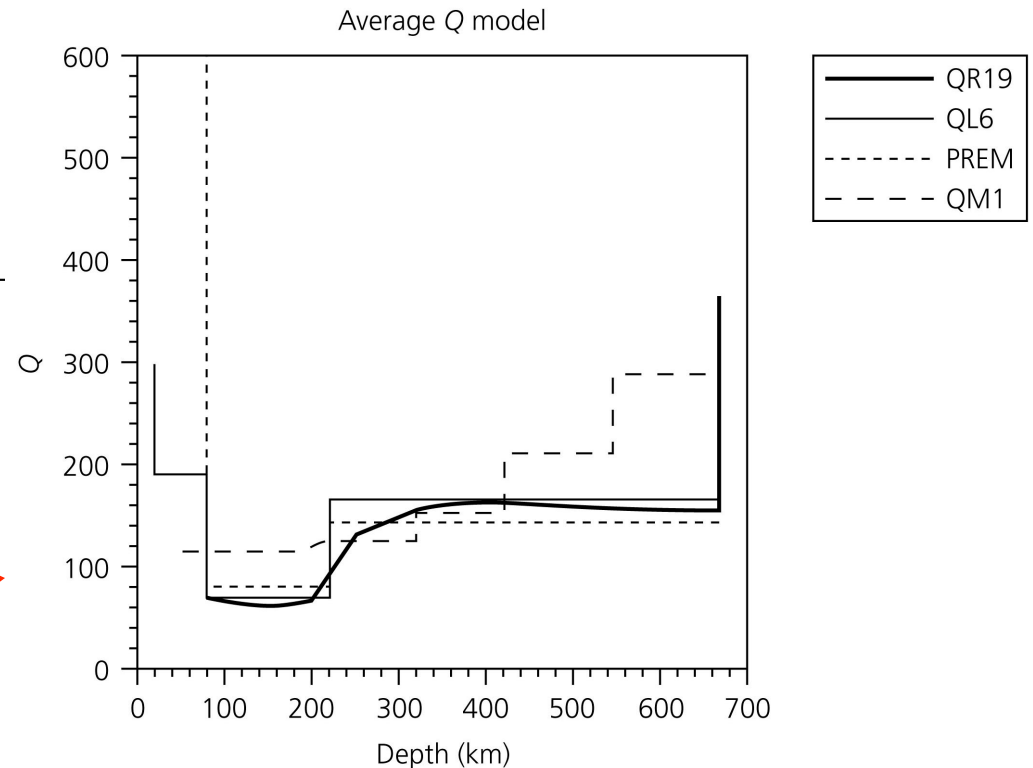


Eastern USA (EUS) and Basin-and-Range Attenuation



Mitchell 1995 (lower attenuation occurring at high frequencies, tells us that frequency independency is not always true

Figure 3.7-19: Models of upper mantle attenuation.



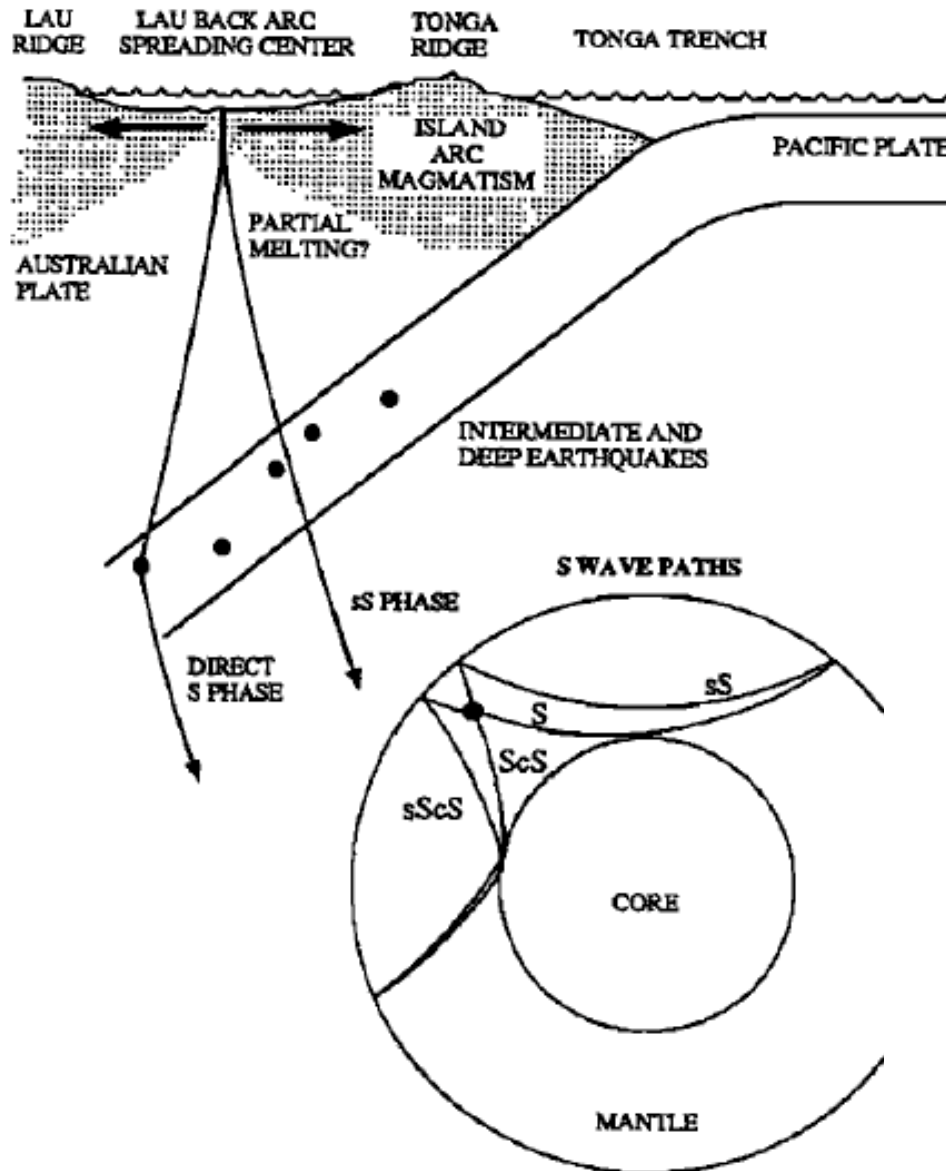
Low Q in the asthenosphere, Romanowicz, 1995

An example of Q extraction from differential waveforms.

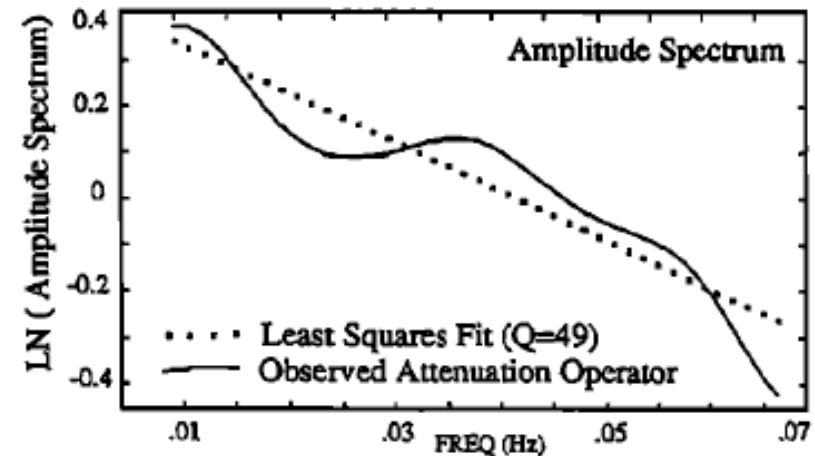
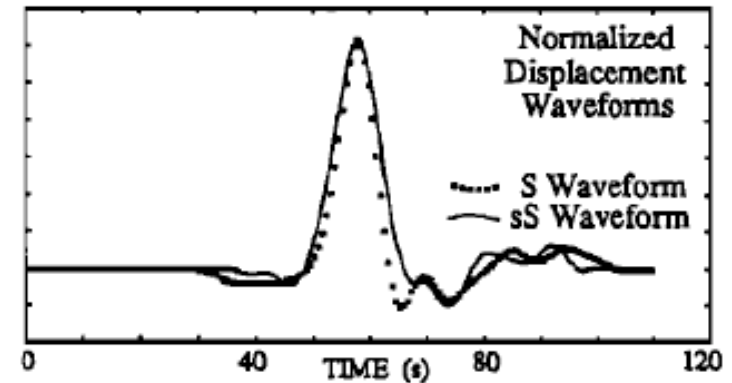
Two phases of interest:

sS-S

sScS - ScS



FREQUENCY DOMAIN ATTENUATION DETERMINATION
APRIL 25, 1984 DEPTH 413 km STATION RSON



Flanagan & Wiens (1990)

Key Realization: The two waveforms are similar in nature, mainly differing by the segment in the above source (the small depth phase segment)

$$sS(\omega) \approx S(\omega)R(\omega)A(\omega)$$

$sS(\omega)$ = frequency spec of sS

$S(\omega)$ = frequency spec of S

$R(\omega)$ = crustal operator

$A(\omega)$ = attenuation operator of interest

Approach: Spectral division of S from sS , then divide out the Crustal operator (a function in freq that accounts for the of the additional propagation in a normal crust)

Spectral dividing S and R will leave the attenuation term $A(\omega)$

$$|A(\omega)| = e^{-\omega t / 2Q} \Rightarrow \log(|A(\omega)|) = -\frac{t}{2Q}\omega \quad \begin{array}{l} \text{(slope} = -t/2Q, \\ t = \text{time diff of } sS-S) \end{array}$$